



Peridynamic Modeling of Material Failure

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Supercomputing at Berkeley when I was a student



CDC 6600: 10MHz



Seymour Cray

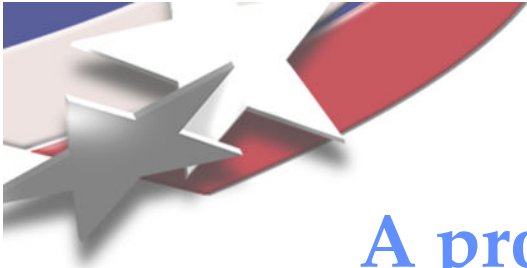


Collaborators

- Prof. Rohan Abeyaratne, Markus Zimmermann, and Olaf Weckner (MIT)
- Abe Askari, John Haws, and Xifeng Xu (Boeing)
- Prof. Kaushik Bhattacharya and Kaushik Dayal (Caltech)
- Prof. Florin Bobaru (University of Nebraska)
- Paul Demmie (Sandia)
- Simon Kahan (Cray)
- Prof. Walter Gerstle (University of New Mexico)
- Kirsten Fagnon (University of Washington)

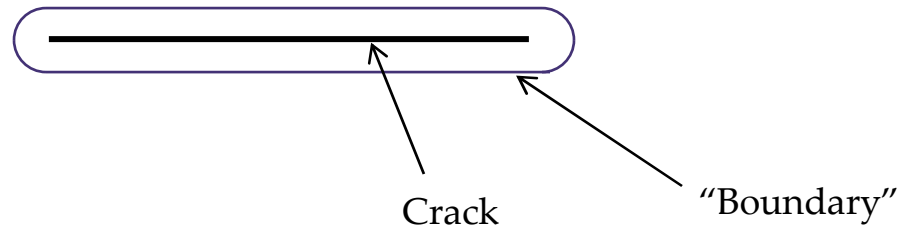
Sponsors

- US Department of Energy
- DoD/DOE Joint Munitions Technology Program
- US Nuclear Regulatory Commission
- The Boeing Company



A problem with the classical theory

- PDEs don't apply when a crack or other discontinuity appears.
 - So cracks have to be treated by special techniques.
 - Example: redefine the body to exclude a crack:



(This doesn't work too well if you don't know where the crack is!)

- Purpose of the peridynamic model:
 - Reformulate the basic equations so that they hold everywhere in a body regardless of discontinuities.



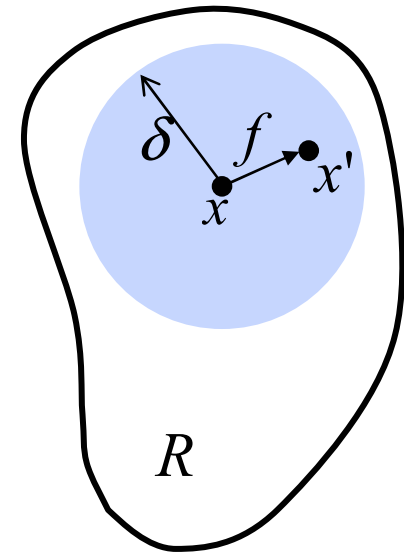
Peridynamic* model

- Replace the $\nabla \cdot \sigma$ term in the equation of motion:

$$\rho \ddot{u}(x, t) = \int_R f(u' - u, x' - x) dV' + b(x, t)$$

- Note the similarity to molecular dynamics.
- f is the force that x' exerts on x per unit volume squared, dependent on:
 - relative position in the reference configuration,
 - relative displacement,
 - (will consider history dependence later).
- **Not** obtainable by applying the divergence theorem to the classical PDE.
- Convenient to assume f vanishes outside some horizon d .
- Require:

$$f(-\eta, -\xi) = -f(\eta, \xi) \quad f(\eta, \xi) \times (\eta + \xi) = 0$$



* From the Greek “near” + “force”



Some references

- Similar idea proposed for multiscale (linear only):
 - I. A. Kunin, *Elastic Media With Microstructure* (1982).
 - D. Rogula, *Nonlocal Theory of Material Media* (1982).
- 3D nonlinear theory:
 - S. A. Silling, Reformulation of elasticity theory for discontinuities and long-range forces, *JMPS* (2000).
 - M. Zimmermann, thesis (to appear).
- Analytical approaches to 1D problems:
 - S.A. Silling, M. Zimmermann, and R. Abeyaratne, Deformation of a peridynamic bar, *Journal of Elasticity* (2003).
 - O. Weckner and R. Abeyaratne, The effect of long-range forces on the dynamics of a bar, *JMPS* (to appear).
- Numerical method:
 - S.A. Silling and E. Askari, A meshfree method based on the peridynamic model of solid mechanics, *Computers and Structures* (to appear).
- Fracture and damage (mostly numerical):
 - S.A. Silling and E. Askari, Peridynamic modeling of impact damage, *ASME PVP-Vol. 489* (2004).
 - S.A. Silling and F. Bobaru, Peridynamic modeling of membranes and fibers, *International Journal of Non-Linear Mechanics* (2005).
- Phase boundaries:
 - K. Dayal, thesis (to appear).



Microelastic materials

- A body is microelastic if f is derivable from a scalar **micropotential** w , i.e.,

$$f(\eta, \xi) = \frac{\partial w}{\partial \eta}(\eta, \xi) \quad \eta = u' - u \quad \xi = x' - x$$

- Interactions (“bonds”) can be thought of as elastic (possibly nonlinear) springs.
- Elastic energy is stored reversibly:

$$\dot{\Phi} = \int_R b \cdot \dot{u} dV$$

- where the strain energy density is

$$W(x) = \frac{1}{2} \int_R w(u' - u, x' - x) dV'$$

- and the total strain energy is

$$\Phi = \int_R W(x) dV$$

- Can show (using Stokes’ theorem) that the force magnitude depends on η only through the current scalar distance between x and x' .

$$w(\eta, \xi) = \hat{w}(|\xi + \eta|, \xi)$$



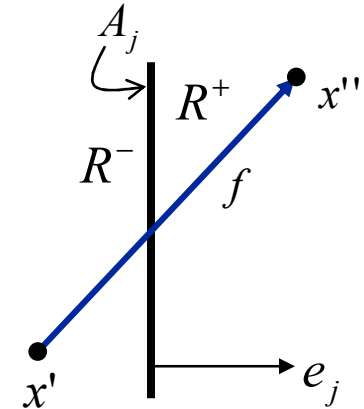
Relation to classical theory

- For a given microelastic material with micropotential w , we can *define* a **classical hyperelastic** material through

$$\hat{W}(F) = \frac{1}{2} \int_R w((F-1)x, x'-x) dV'$$

- Can *define* a stress-like quantity

$$\sigma_{ij}(x) = \lim_{A_j \rightarrow 0} \left\{ \frac{1}{A_j} \int_{R^+} \int_{R^-} f_i(u''-u', x''-x') dV'' dV' \right\}$$



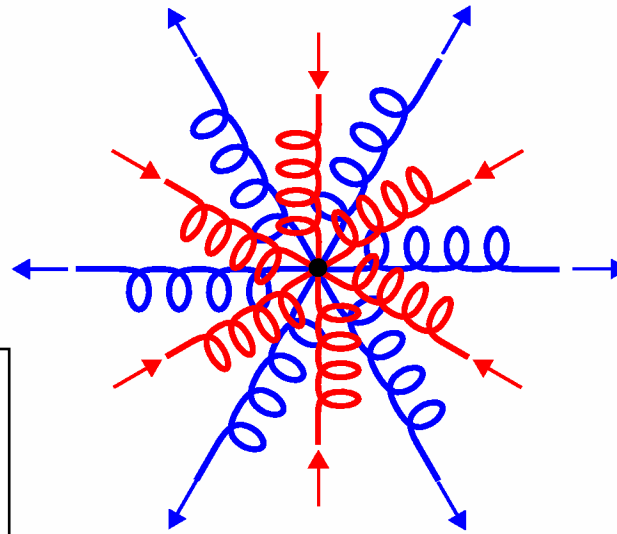
but this is meaningful only for homogeneous deformations.

- Can show that the peridynamic equation of motion “converges to” the classical version in the limit $\delta \rightarrow 0$.



Unstressed configurations

- Forces between particles can be nonzero even when the “stress” vanishes.
- Could make this an attractive approach for “multiscale” modeling.



Unstressed: Bond forces are not only equilibrated, but their sum through any surface vanishes.

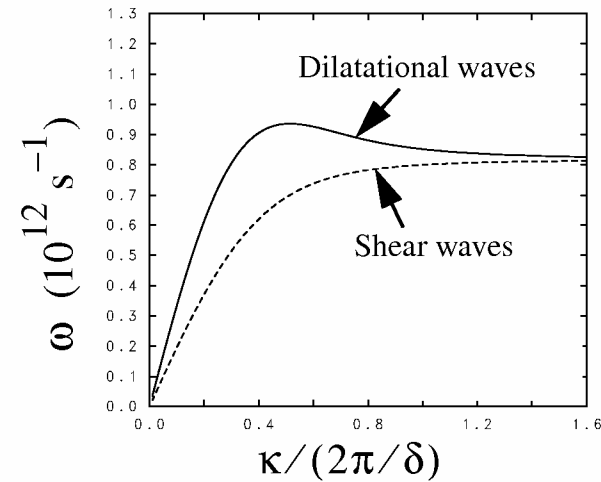


Relation to classical theory, ctd.

- Can show that an unbounded, homogeneous, isotropic peridynamic body sustains two types of small-amplitude waves:

$$\underline{u}(\underline{x}, t) = \underline{a} e^{i(\underline{\kappa} \cdot \underline{x} - \omega t)}$$

- Longitudinal (displacement parallel to propagation direction)
- Shear (displacement orthogonal to propagation direction)
- But these waves are dispersive.



Typical dispersion curves

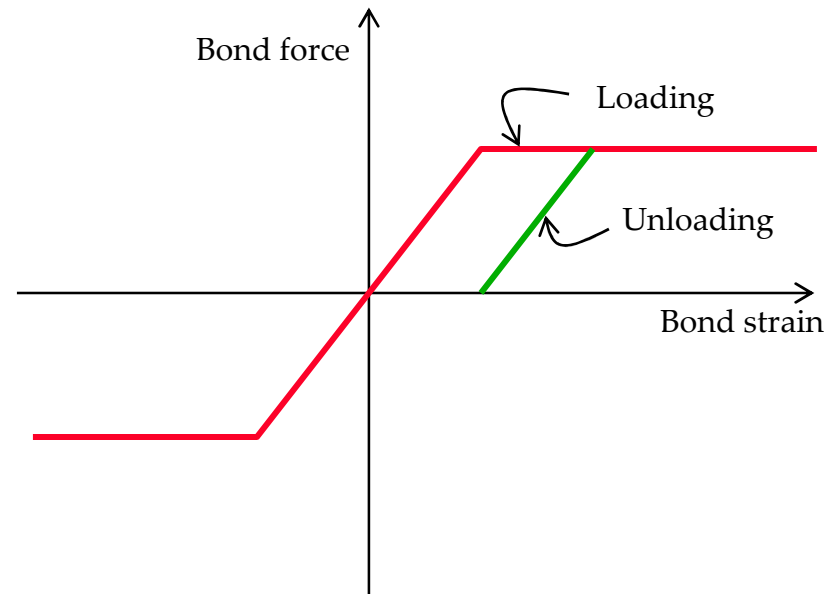


Other constitutive models

- Visco-microelastic:

$$f(r, \dot{r}, \xi), \quad r = |\xi + \eta| = \text{current distance between } x \text{ and } x'$$

- Microplastic:





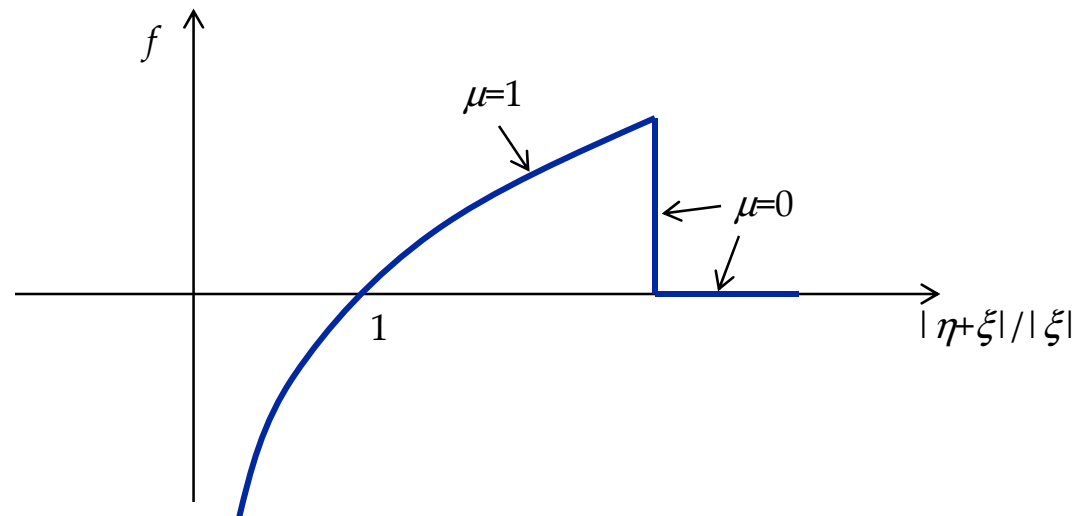
Damage

- Damage is introduced at the bond level:

$$\bar{f}(\eta, \xi, x, t) = f(\eta, \xi) \mu(\xi, x, t)$$

where $\mu = 1$ for an intact bond, 0 for a broken bond.

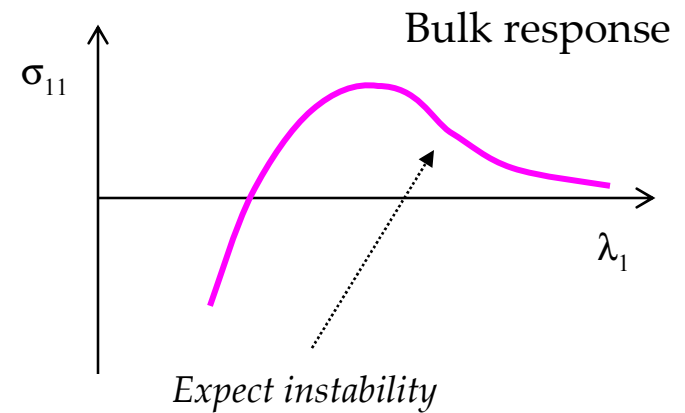
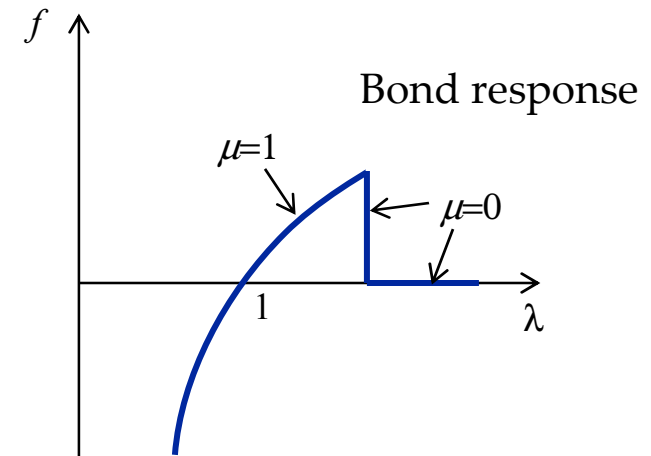
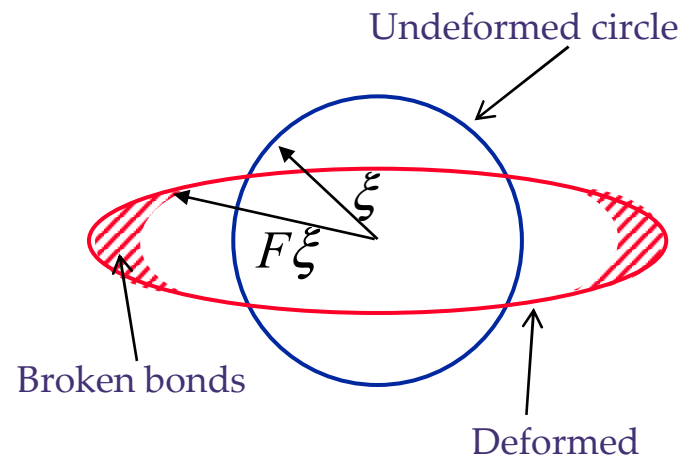
- Bond breakage occurs irreversibly according to some criterion such as exceeding a prescribed critical stretch.





Bulk response with damage

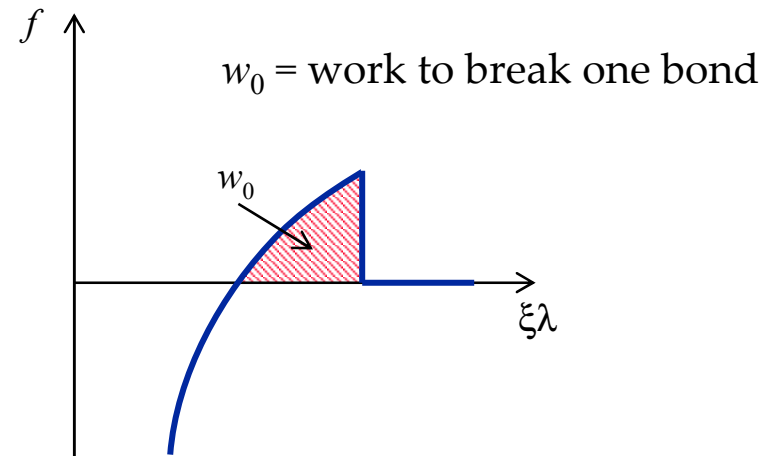
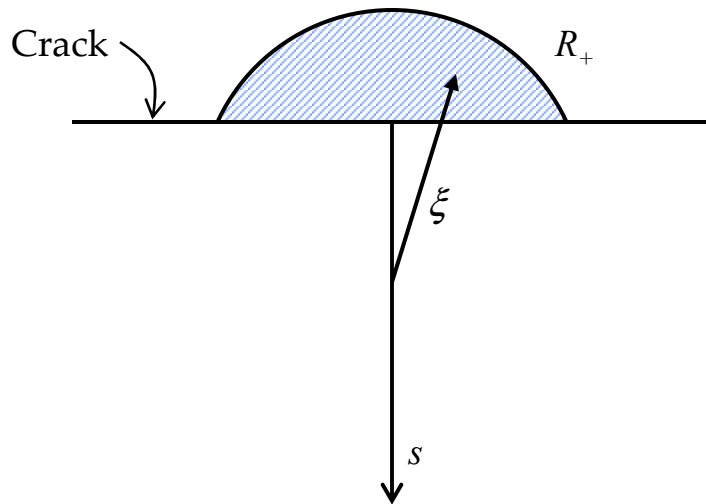
- Assume a homogeneous deformation.



Energy required to advance a crack

- Adding up the work needed to break all bonds across a line yields the energy release rate:

$$G = 2h \int_0^\delta \int_{R_+} w_0 dV ds$$



There is also a version of the J-integral that applies in this theory.



Numerical method

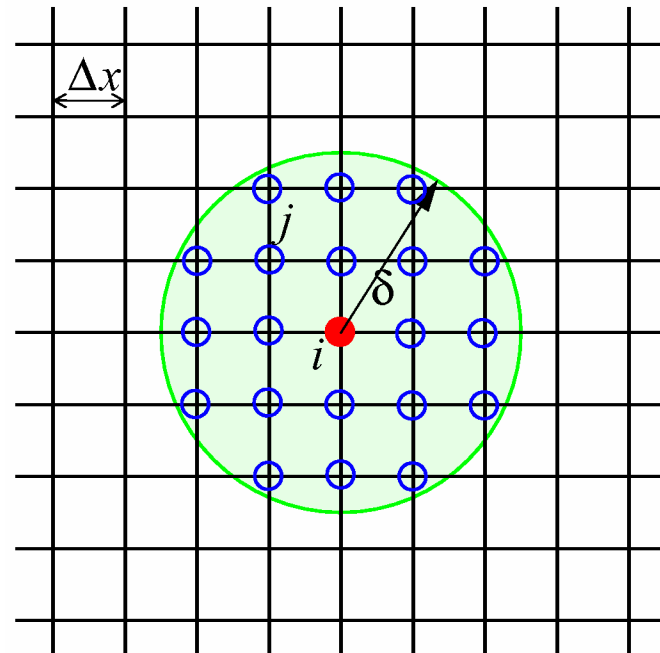
- Replace the integral in the equation of motion by a finite sum:

$$\rho \ddot{u}(x, t) = \int_R f(u' - u, x' - x) dV' + b(x, t)$$

is approximated by

$$\frac{\rho}{\Delta t^2} (u_i^{n+1} - 2u_i^n + u_i^{n-1}) = \sum_{|x_j - x_i| < \delta} f(u_j^n - u_i^n, x_j - x_i) (\Delta x)^3 + b_i^n$$

- Method is “meshless Lagrangian”.
 - no elements
 - error is $O(\Delta x^2)$ if u is continuous.
- Sandia 3D peridynamic code is called Emu.
 - available under license from Sandia

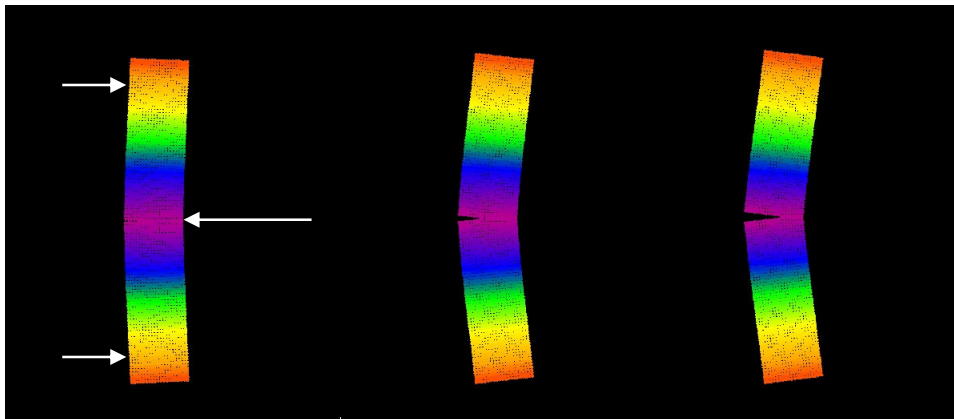


Special case of a cubic grid

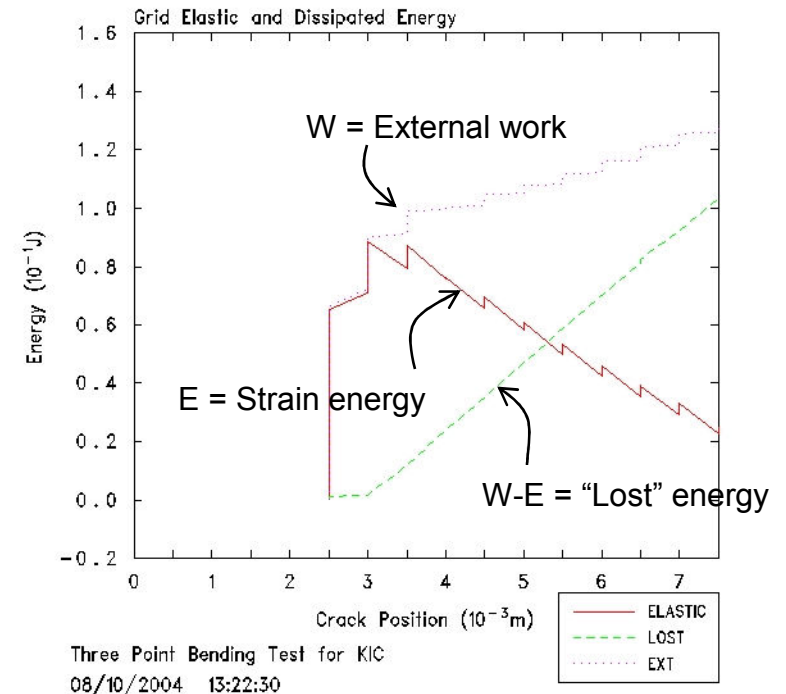


3-point bend test in a metal: Emu model

- 2D model of ASTM standard test for K_{Ic}.



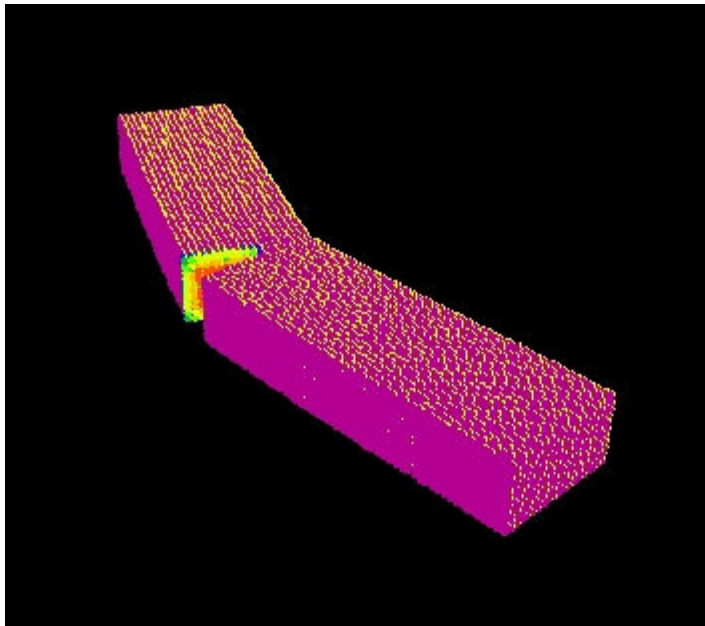
- Constant slope of the “lost” energy curve is the energy release rate G .
- Confirms that the code accurately models fracture under the assumption of constant K_{Ic}.
- Also shows how this quantity depends on model parameters.



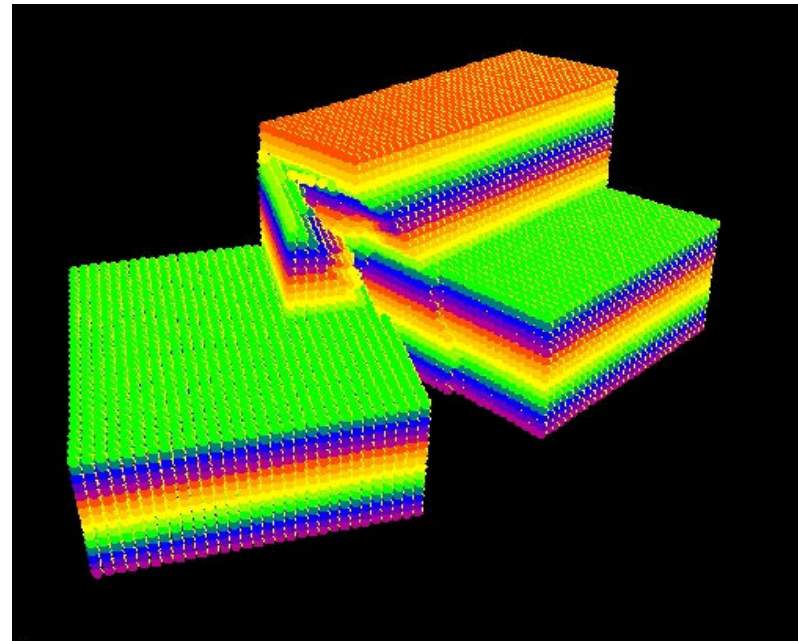
Energy as a function of crack length



Single crack growth in metals: 3D model



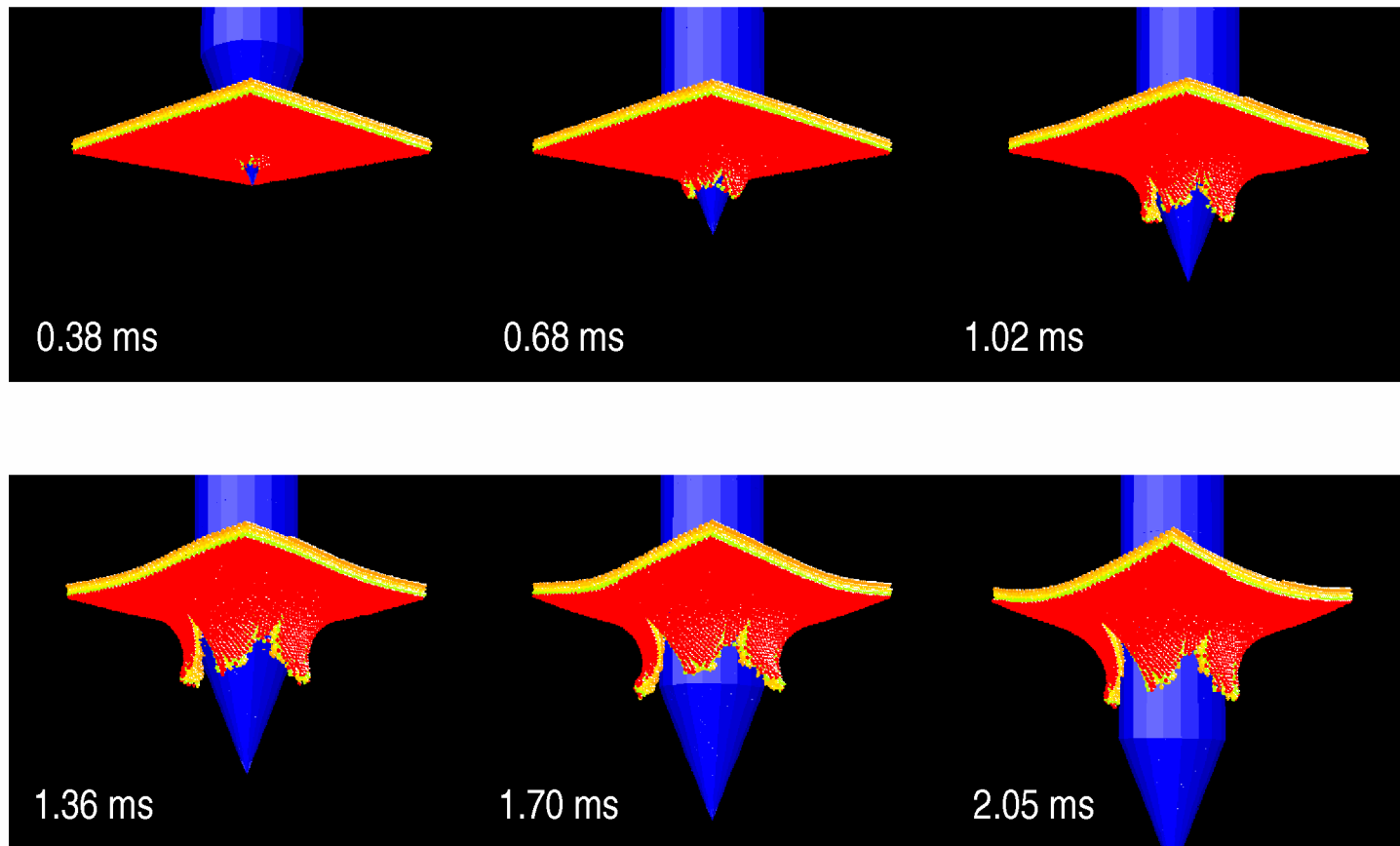
3-point bend



A more complex geometry

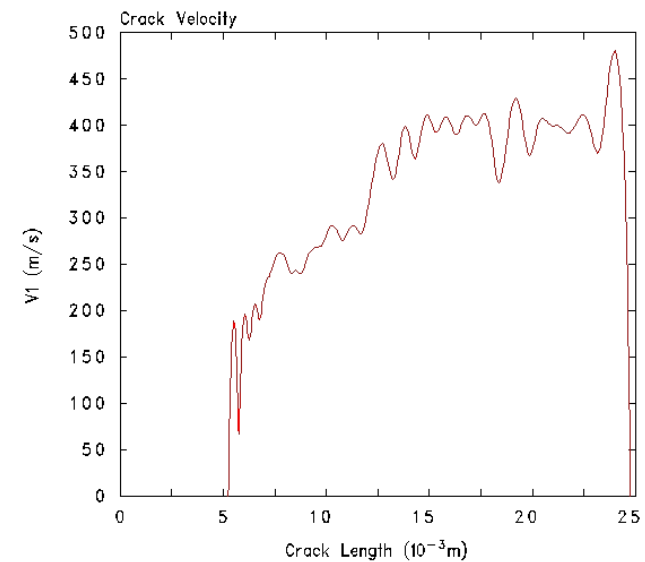
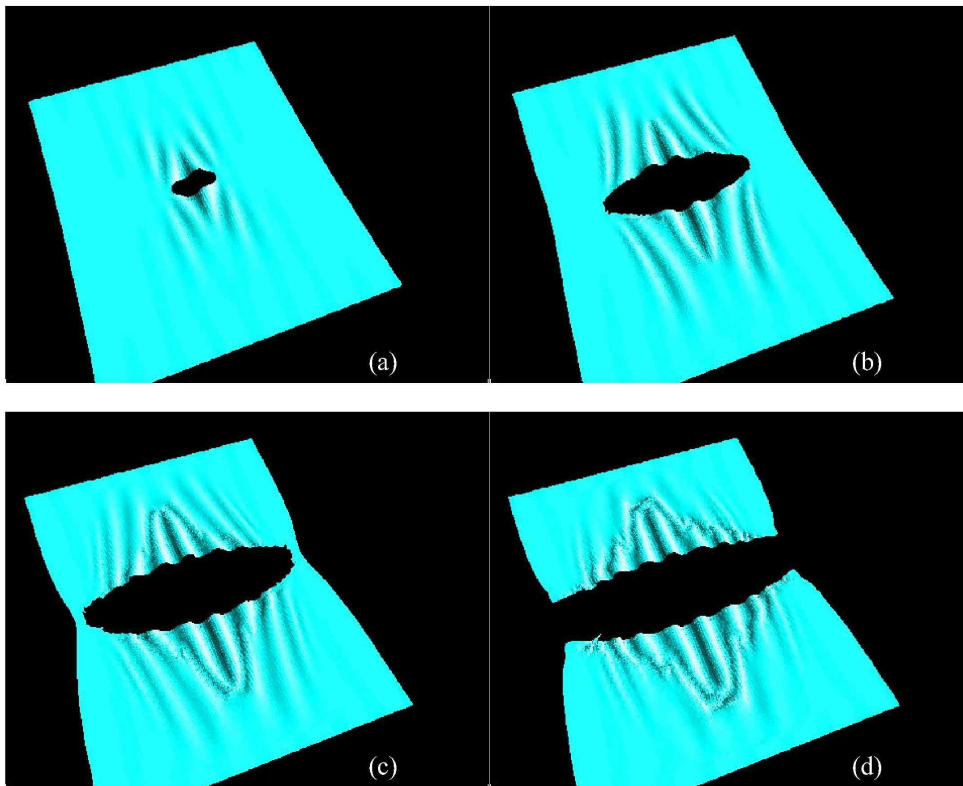


Petaling in perforation of a ductile plate



Dynamic fracture: Tearing of a membrane

- Wrinkles appear due to compressive strains parallel to the crack*.



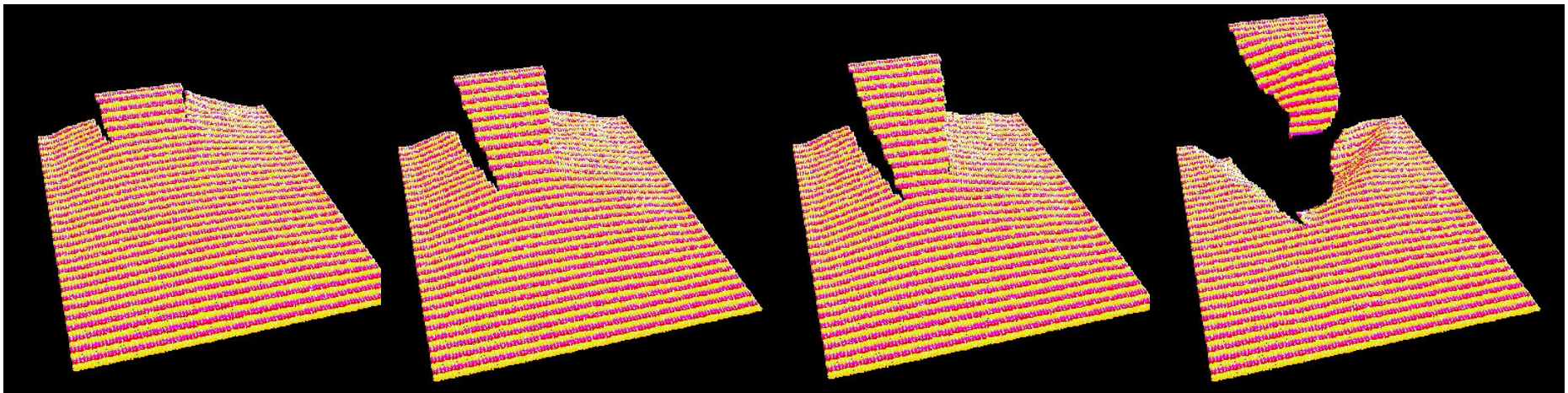
Crack tip velocity

*Also see Haseganu and Steigmann, *Computational Mechanics* (1994) for numerical model of wrinkling.

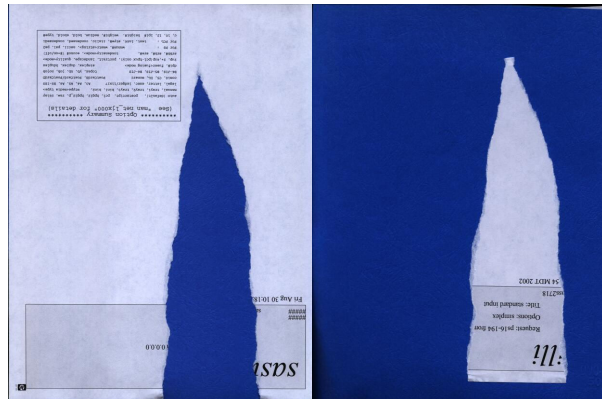


Interaction of 2 cracks: Peeling of a sheet

- Pull upward on part of a free edge – other 3 edges are fixed.

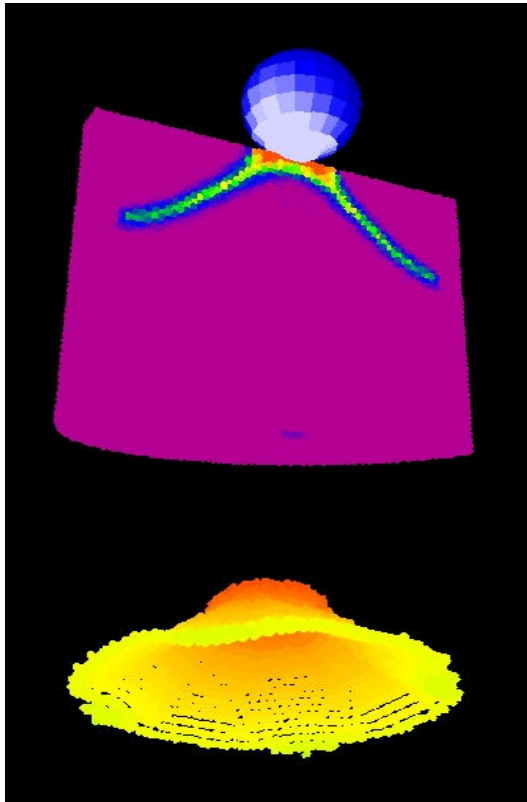


“Experimental data”

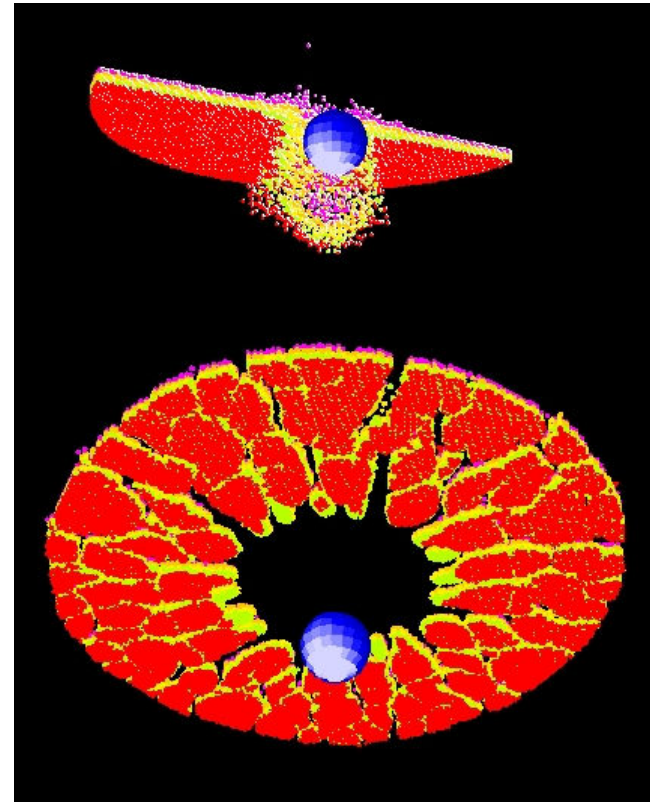




Hertzian cracking and fragmentation in glass



Low speed impact forms
conical crack

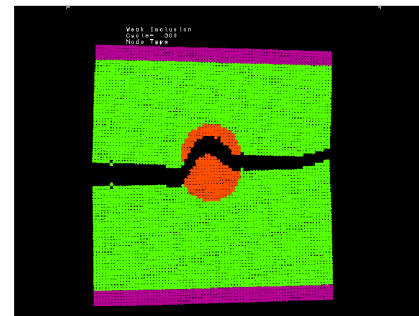
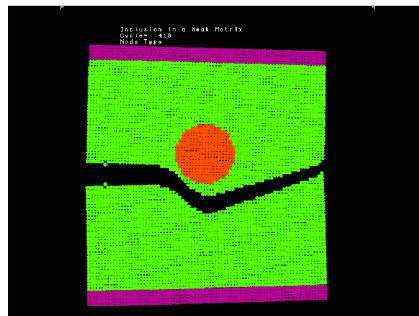
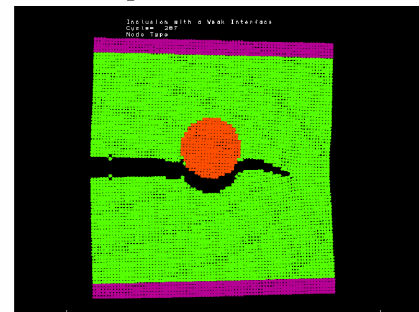
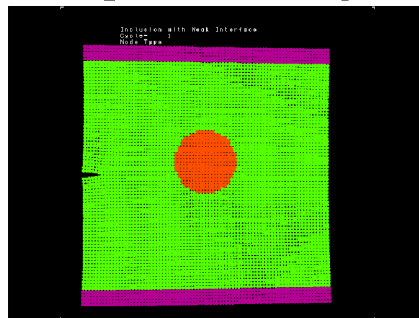


Higher speed impact forms
fragments



Treatment of interfaces

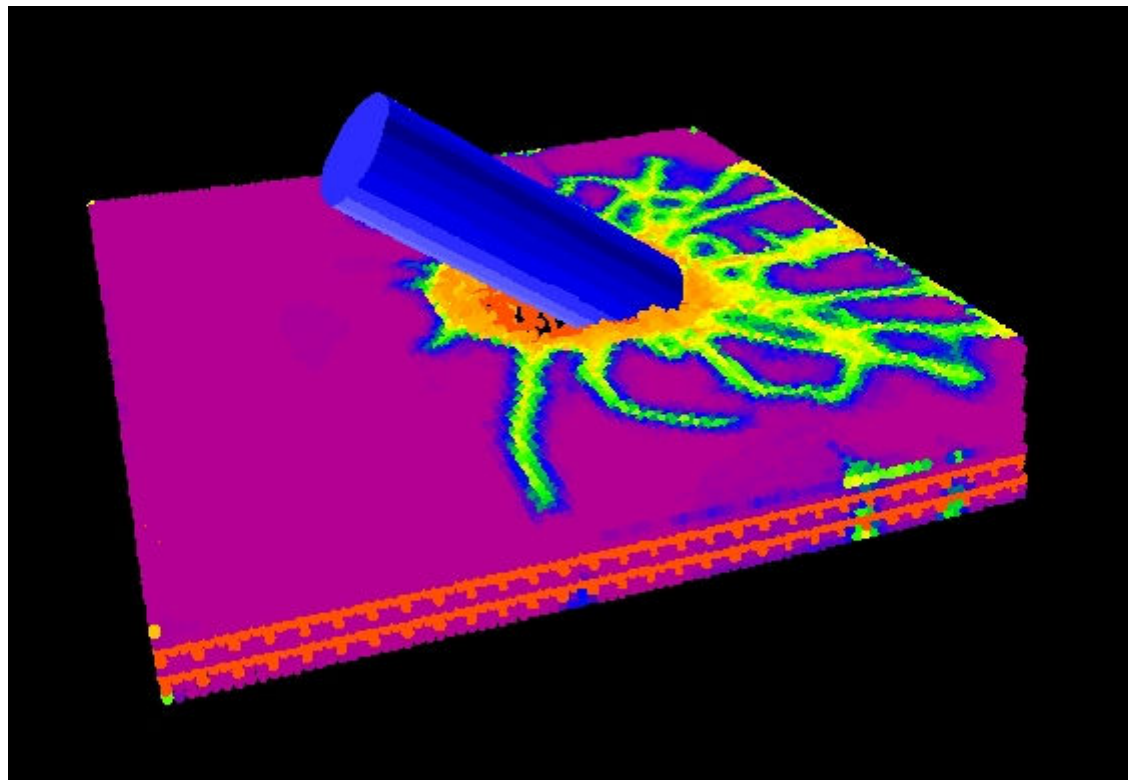
- Bonds connecting different materials can have properties independent of the constituent properties.
- Crack growth is “autonomous:” no need for supplemental kinetic relations.





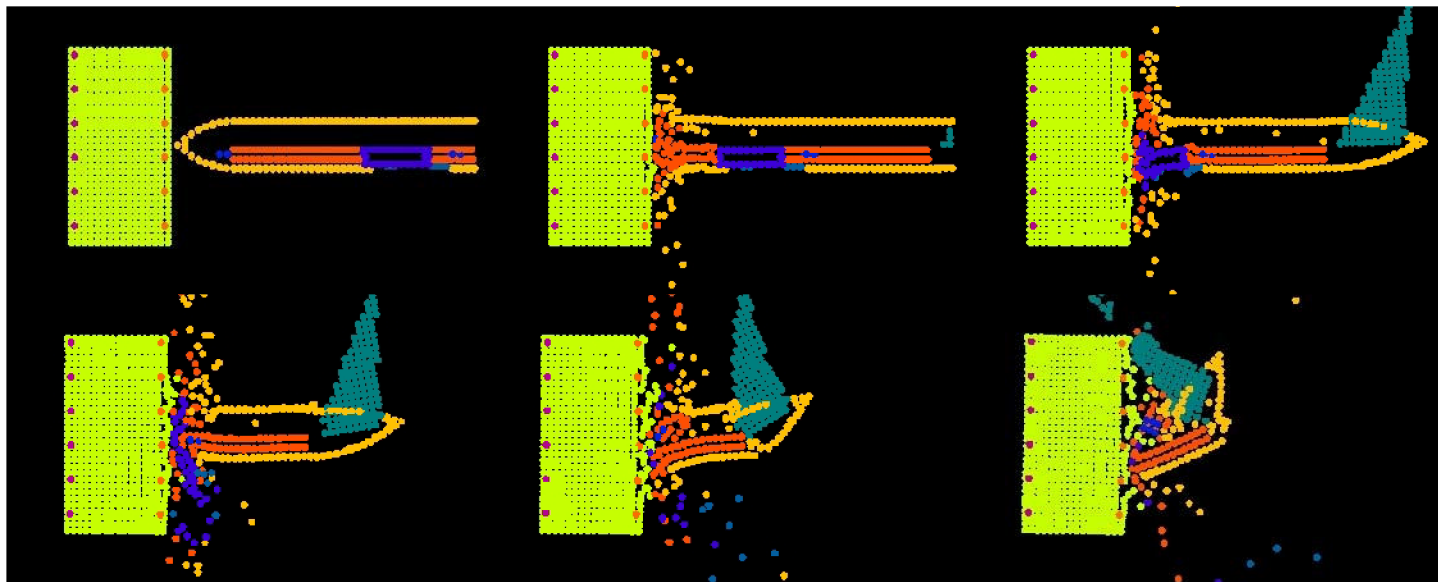
Impact onto reinforced concrete

- Reinforcement is modeled explicitly.





F4 airplane into concrete block

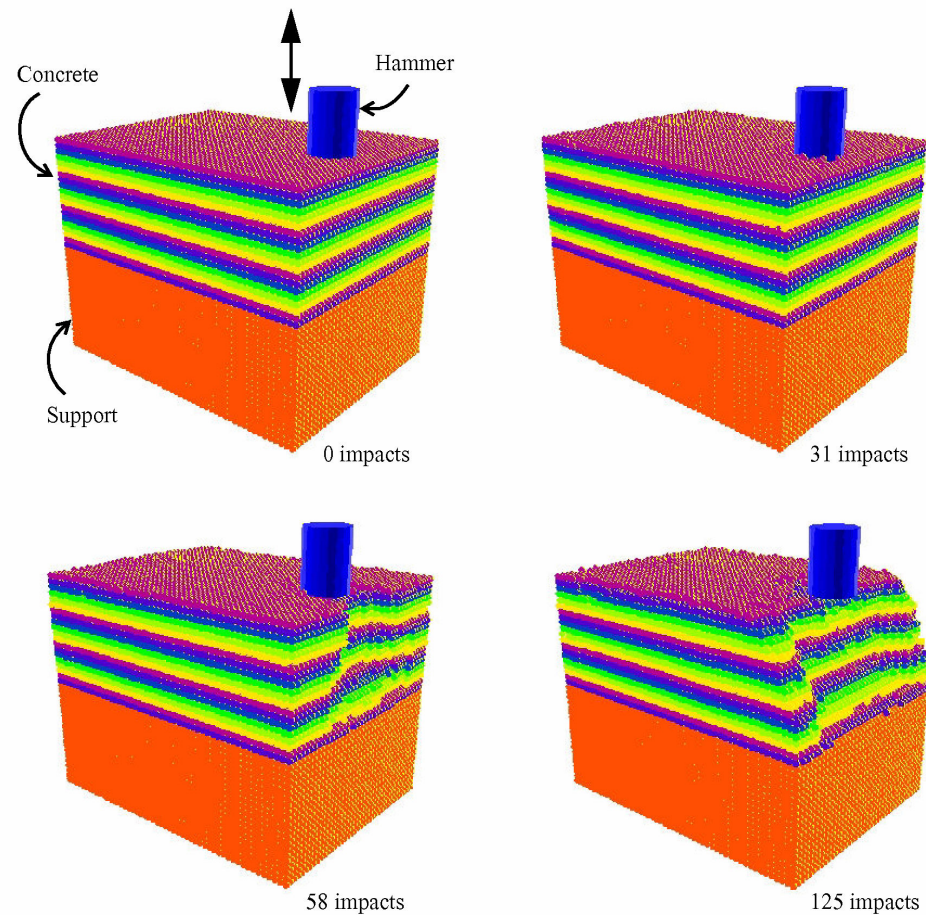


- EMU model of full-scale experiment (Sugano et al, Nuclear Engineering and Design 140 373-385 (1993)).



Damage accumulation due to repetitive impact (“jackhammer”)

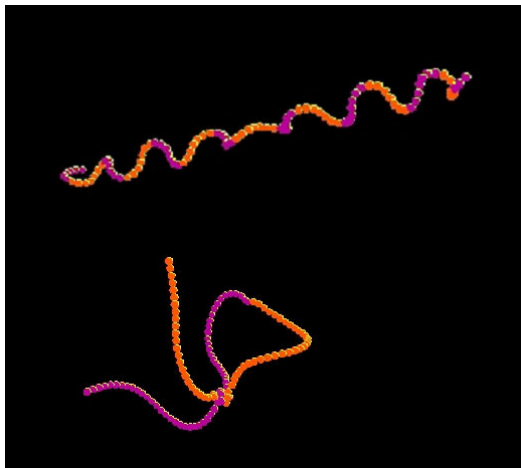
- Each successive impact breaks more bonds internally.
- These coalesce into large cracks.



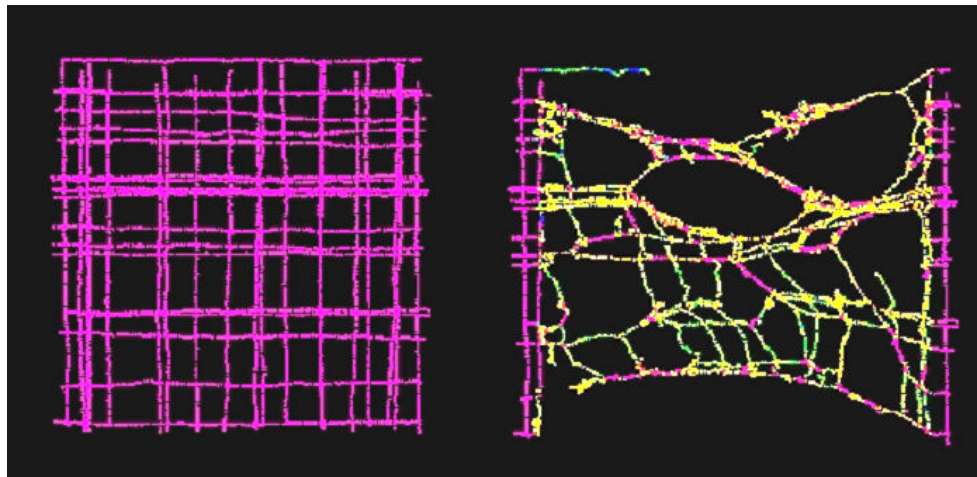
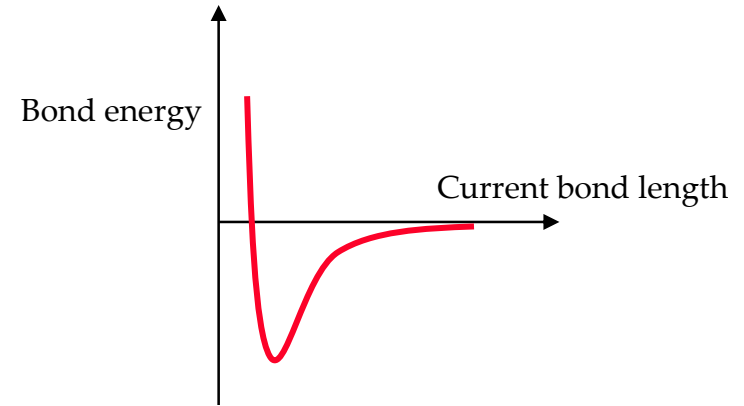


Fibers and fiber networks

- Long-range (e.g. van der Waals) forces treated the same way as peridynamic forces.



Fibers in which different segments attract or repel each other

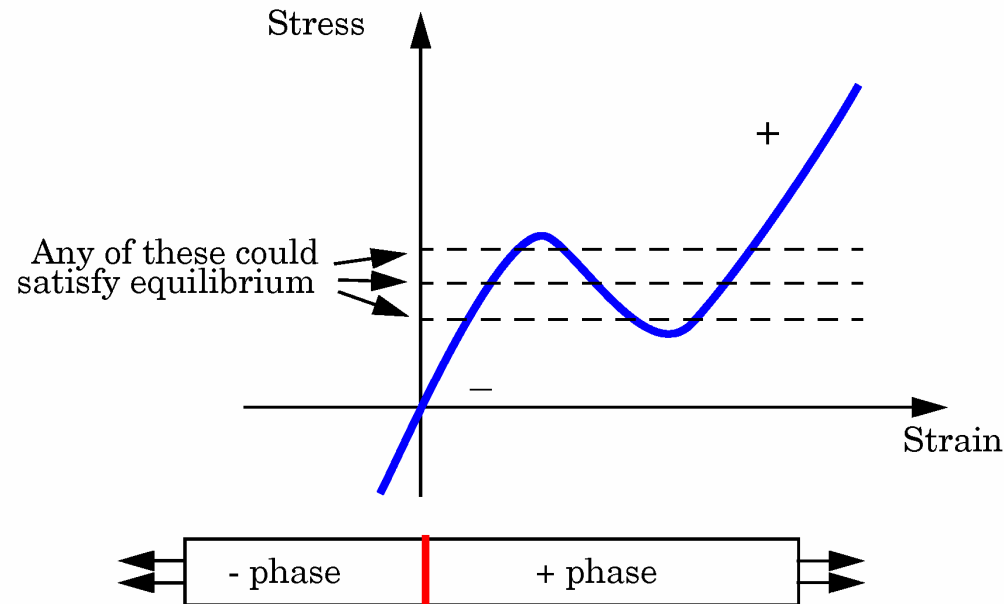


Failure of a nanofiber network including VDW forces
(F. Bobaru, Univ. of Nebraska)



Phase boundaries: classical approach

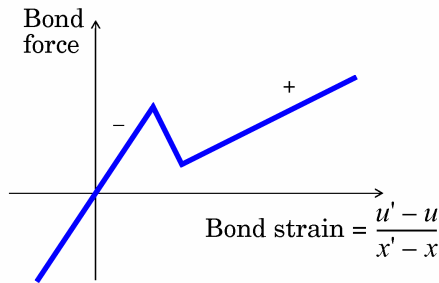
- Classical theory:
 - Supplemental condition (such as the Maxwell condition) is required to determine the conditions at a phase boundary.



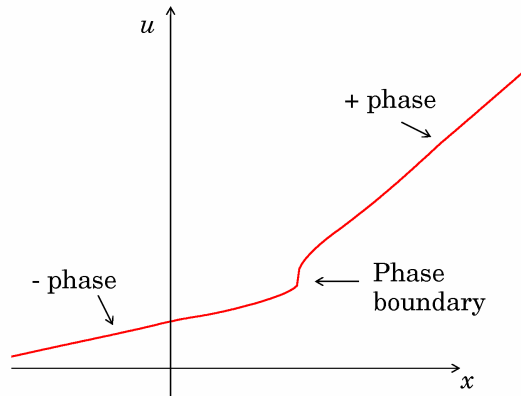


Phase boundaries: peridynamic

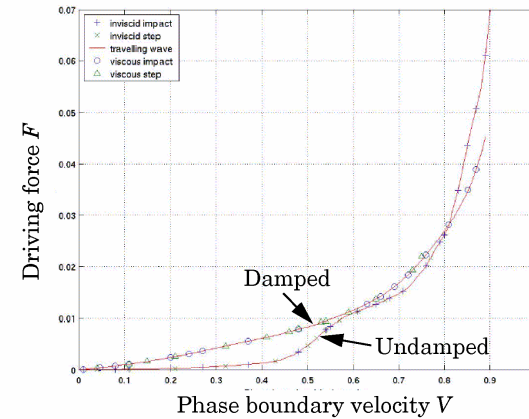
- Constitutive model:



- Numerical solution for static displacement field:



- Computed phase boundary velocity (for a specific material):



where the driving force (in the classical theory) is defined by

$$F = [W] - \langle \sigma \rangle [\varepsilon].$$

Peridynamic model appears to select particular conditions across the phase boundary (results of K. Dayal).

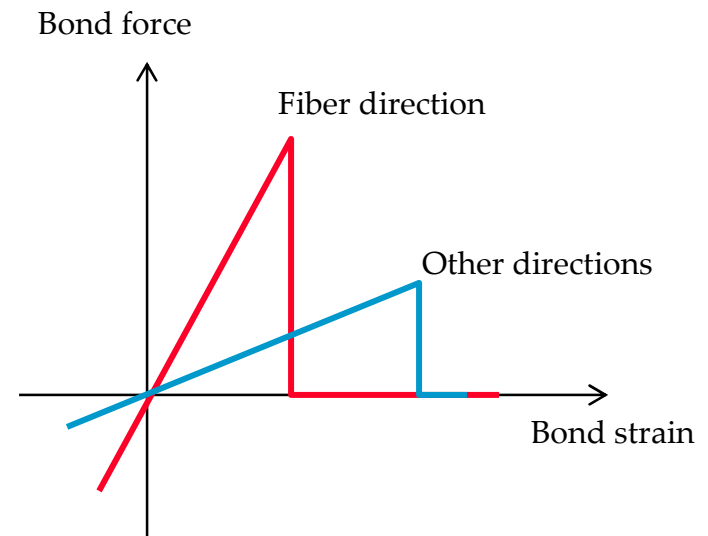
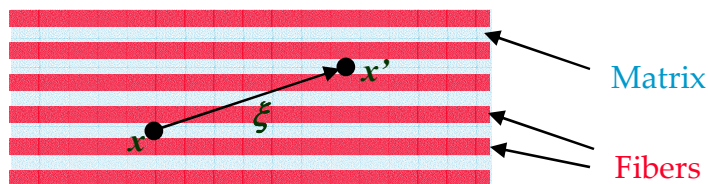


Laminated composite materials

- Recall that the bond force can depend on the separation vector ξ between x and x' in the reference configuration.

$$f(\eta, \xi)$$

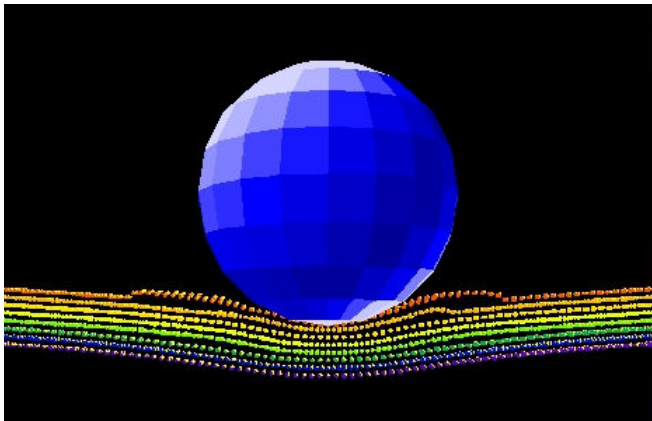
- Bonds in the fiber direction are stiffer and stronger than the others.
 - Micromodulus of each is fitted to bulk elastic modulus in each direction.
 - Interface bonds (connecting 2 different materials) can have properties independent of the others.



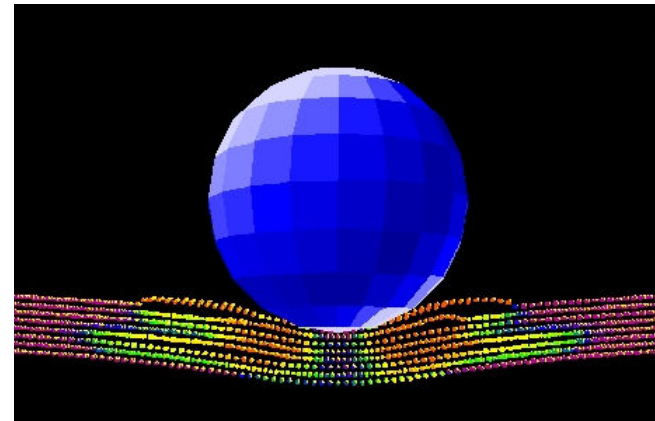


Impact on a laminated composite

- 25/50/25 stacking sequence, 32 layers, 1" diameter projectile at 5.1 m/s).
- Note blistering near front surface due to delamination + buckling.
- Central area has less damage because of large out-of-plane compressive stress.
 - Also the shear stress vanishes along the central axis.



Colors indicate material (layer).



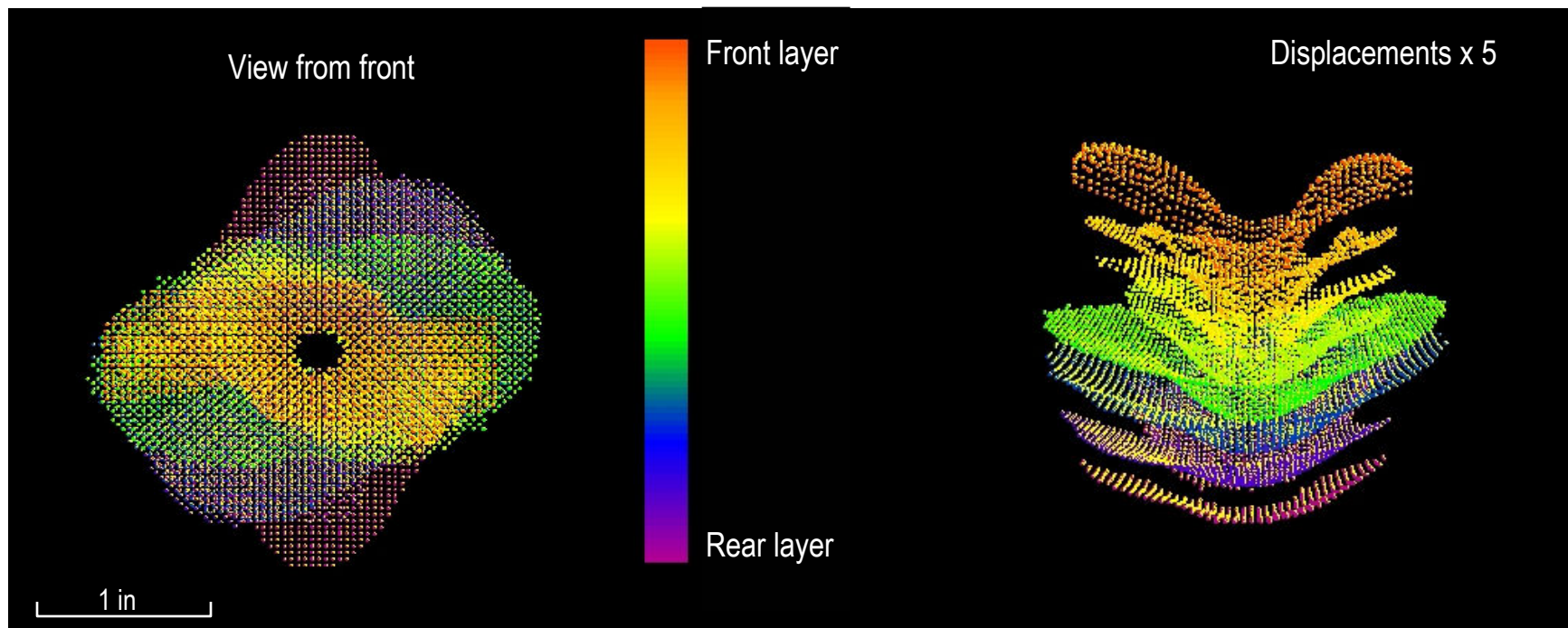
Colors indicate damage.

Both figures are at time of maximum penetration depth



Impact on a composite: Views of delaminated areas

- Delaminations form in roughly elliptical regions.

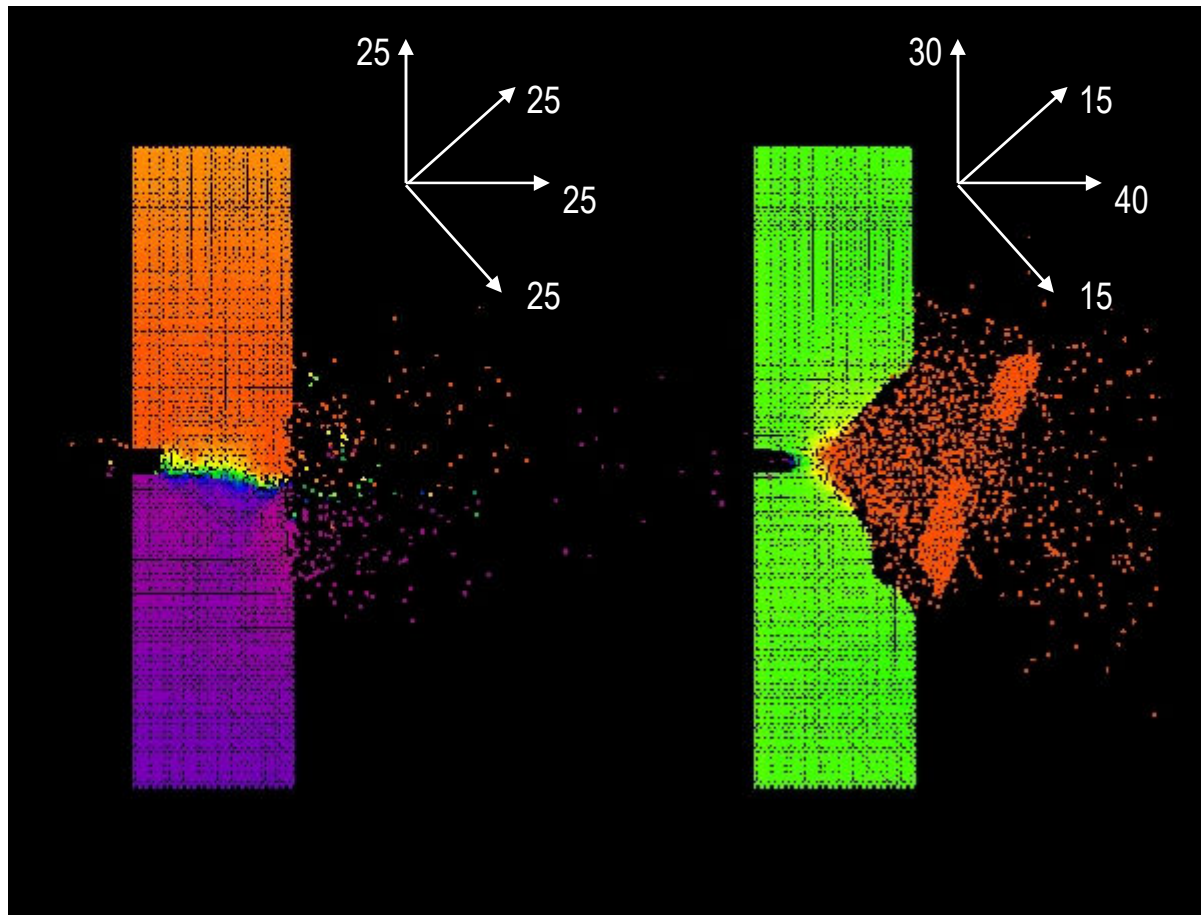


Colors indicate material (layer)
Only damaged layers are shown.



Fracture modes in a laminated composite

- How does the makeup of a composite influence how it fails in a middle tension test (with a blunt notch)?





Some research issues

- Criteria for material stability and fracture:
 - Convexity of some quantity related to the micropotential?
- Plasticity with plastic incompressibility.
- Quantitative relationship with molecular dynamics.
 - Multiscale: method for “coarse graining”.
 - Nanoscale.
- Homogenization of heterogeneous materials.

For further information:

- www.sandia.gov/emu/emu.htm
- Email to sasilli@sandia.gov